

## Electronic and Microscopic Properties of Spin $\frac{1}{2}$ System, $\text{CuInO}(\text{VO}_4)$

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### Abstract

We present first principles density functional theory (DFT) calculations within the generalized gradient approximation for the exchange-correlation functional in order to gain microscopic understanding, electronic and microscopic properties of  $\text{CuInO}(\text{VO}_4)$ . We discuss explicitly the electronic properties, nature of the exchange paths and estimate magnetic exchange interactions. Focusing on the calculated magnetic interactions and electronic structure of  $\text{CuInO}(\text{VO}_4)$  puts it in the class of weakly coupled low dimensional antiferromagnetic tetramer with spin  $S=1/2$ .

**Keywords:** Electronic Structure, Magnetic Exchange Interaction, Wannier Function, Microscopic Modelling

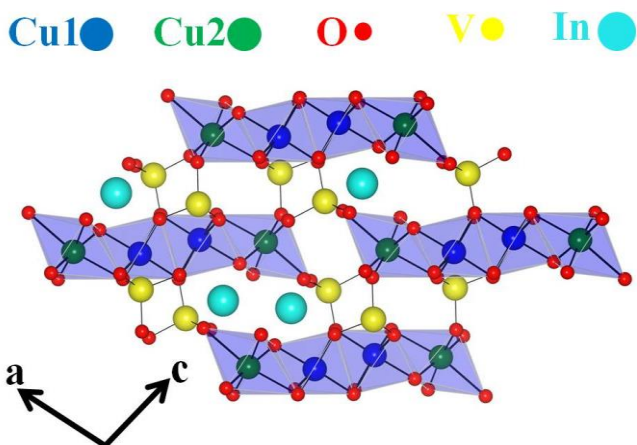
### 1. Introduction

The importance of low dimensional quantum spin systems is due to quantum fluctuation in the low energy regime. Spin fluctuation is large for small spin and low dimensional systems. Odd- and even-leg ladders, dimer, alternating chain, spin-peierls are the different class of low dimensional quantum spin systems. Out of these dimer[1], spin  $\frac{1}{2}$  or spin 1 alternating Heisenberg antiferromagnetic (AFM) chains[2,3], even-leg ladders[4], spin-peierls[5] are spin gap systems whereas odd leg ladders[6], spin  $\frac{1}{2}$  uniform Heisenberg AFM chain[2,3] are spin gapless systems. Like dimer, tetramer chain systems have a finite gap between singlet ground state and triplet excited state.  $\text{CuInO}(\text{VO}_4)$  is the class of low dimensional quantum spin system with spin  $S=1/2$ . The first-principles modelling of materials, which takes into account the chemistry and structural information accurately, turns out to be a useful tool in this context. This approach has been immensely successful in the modelling of various different spin compounds[7]. In this study, considering the case of  $\text{CuInO}(\text{VO}_4)$ , through first-principles calculations we establish that  $\text{CuInO}(\text{VO}_4)$  behaves as weakly coupled spin tetramer with alternating nearest-neighbour.

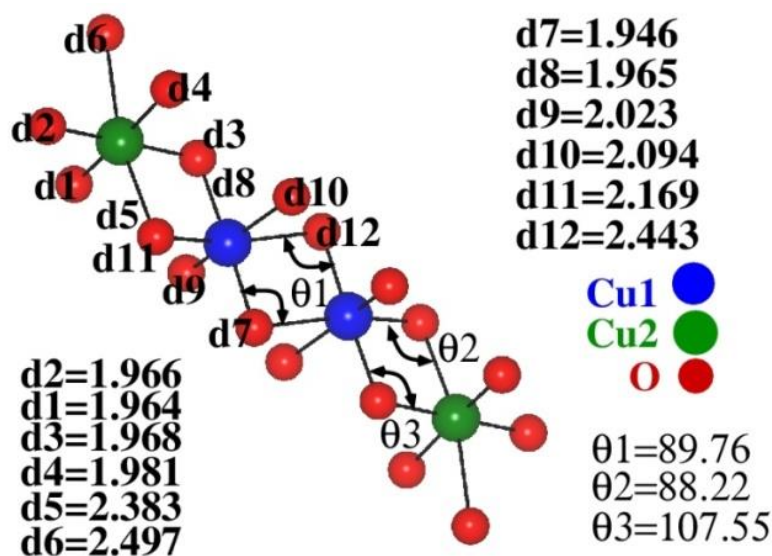
### 2. Crystal Structure

$\text{CuInO}(\text{VO}_4)$  has monoclinic symmetry with space group  $P2_1/c$  (No. 14). Unit cell of  $\text{CuInO}(\text{VO}_4)$  consisting of 8 formula units. There are two inequivalent Cu atoms, Cu1 and Cu2 which are at the centre of a distorted octahedra constructed by six O-atoms. The lattice parameters of the compound are  $a=8.793\text{\AA}$ ,  $b=6.1542\text{\AA}$  and  $c=15.262\text{\AA}$ ,  $\beta=106.69^\circ$ . [8] In  $\text{CuInO}(\text{VO}_4)$ ,  $\text{Cu}_1\text{O}_6$  and  $\text{Cu}_2\text{O}_6$  octahedra share a common edge and form a tetramer  $\text{Cu}_2\text{-Cu}_1\text{-Cu}_1\text{-Cu}_2$  chain. The tetramers are connected by  $\text{VO}_4$  tetrahedra. The voids and channels are occupied by In atoms. The different atoms are represented by

different colours as shown in the Figure 1. The various Cu-O distances (in Å) and angles (in degree) in Cu1O6 & Cu2O6 distorted octahedral of CuInO(VO<sub>4</sub>) is shown in Figure 2.



**Figure 1:** The crystal structure of CuInO(VO<sub>4</sub>) in distorted octahedral representation in ac-plane. The octahedral are connected by edge sharing.



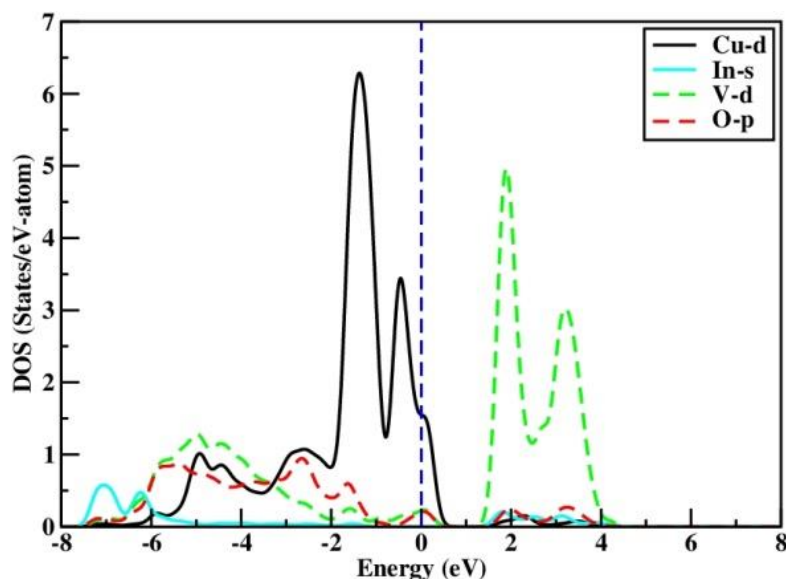
**Figure 2:** The various Cu-O distances (in Å) and angles (in degree) in Cu1O6 & Cu2O6 distorted octahedral in CuInO(VO<sub>4</sub>).

### 3. Electronic Structure

We used first principles density functional theory (DFT)[9] calculations within the GGA [10] for the exchange-correlation functional to investigate the microscopic understanding of CuInO(VO<sub>4</sub>). We have

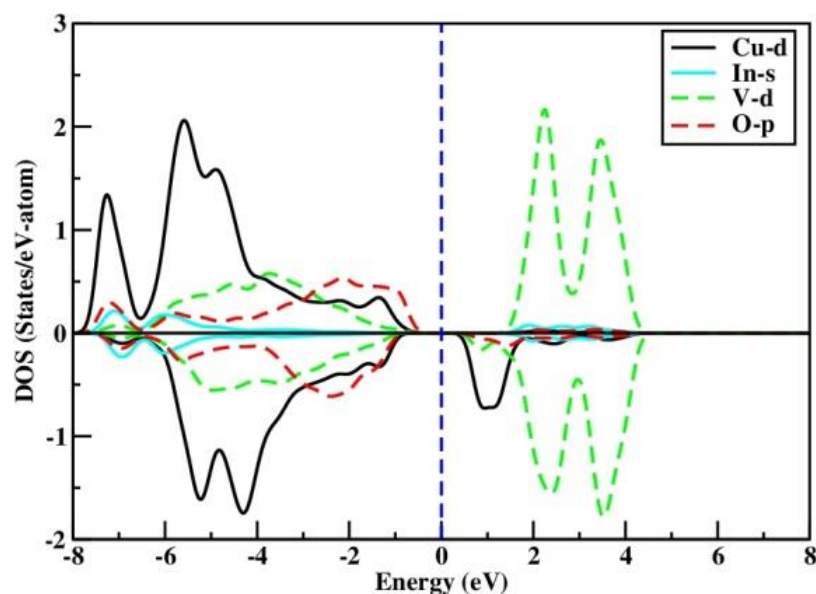
used Vienna An initio Simulation Package (VASP)[11] based on plane wave basis as well  $N$ -th order MTO (NMTO) [12] and linear MTO (LMTO) [13] which are based on muffin-tin orbital (MTO).

The non-spin polarized density of state (DOS) of the compound  $\text{CuInO}(\text{VO}_4)$  is shown in Figure 3. The black colour at Fermi level in the DOS represents  $\text{Cu-d}_{x^2-y^2}$  which is pronouncedly mixed with O-p and V-d states. Therefore, the electronic and magnetic properties of the title compound are mainly depending on the  $\text{Cu-d}_{x^2-y^2}$  state.



**Figure 3:** Electronic density of states (DOS) within GGA of exchange-correlation has been calculated using density functional theory of  $\text{CuInO}(\text{VO}_4)$ . Energy axis is plotted with respect to Fermi energy.

The corresponding spin-polarized density of state is shown in Figure 4 where Cu-d, O-p, V-d and In-s states are represented by black, red, green and cyan colours respectively and energy axis is plotted with respect to Fermi energy. All the d-states of Cu are completely filled in the majority and minority spin-channels except the minority channel of  $\text{Cu-d}_{x^2-y^2}$  which is half filled due to the crystal-field splitting of the distorted octahedra. Therefore, Cu has the nominal  $\text{Cu}^{2+}$  or  $d^9$  valence. The O has also nominal  $\text{O}^{2-}$  valence states because O-p state is found to be mostly occupied. The finite hybridization of O-p and V-d states with Cu-d states close to Fermi level suggesting that the super-exchange path of magnetic interaction between two Cu sites are mediated by O and V atoms. The oxidation state of In and V ions are  $3+$  and  $5+$  respectively. The magnetic moment at Cu, O and V sites are found to be  $0.72 \mu_B$ ,  $0.18 \mu_B$  and  $0.1 \mu_B$  respectively with a total magnetic moment of  $1 \mu_B$  per formula unit.



**Figure 4:** The spin polarized density of state for  $\text{CuInO}(\text{VO}_4)$ . The up channel of Cu-d is completely filled while the down channel of Cu-d is partially filled. Energy axis is plotted with respect to Fermi energy.

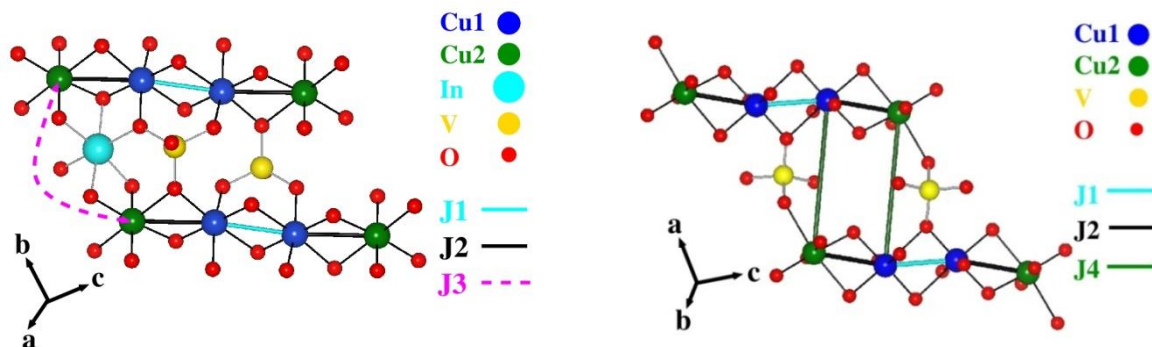
#### 4. Magnetic Interactions

We have used Nth order Muffin Tin Orbital (NMTO) based downfolding technique to find out the various Cu-Cu magnetic exchanges in the compound. Using this technique,  $\text{Cu-d}_{x^2-y^2}$  only wannier function is constructed by keeping active only the  $\text{Cu-d}_{x^2-y^2}$  degrees of freedom and downfolding all the degrees of freedom associated with O, V, In and the Cu. This procedure gives renormalization of  $\text{Cu-d}_{x^2-y^2}$  orbital due to hybridization from O-p, V-d and In-s and the other Cu-d orbitals. The real space representation of the Hamiltonian in the effective  $\text{Cu-d}_{x^2-y^2}$  wannier function basis provides the effective  $\text{Cu-d}_{x^2-y^2}$ – $\text{Cu-d}_{x^2-y^2}$  hopping interactions. Figure 6 represents the effective  $\text{Cu-d}_{x^2-y^2}$  wannier function of  $\text{CuInO}(\text{VO}_4)$ . The central part is shaped according to  $d_{x^2-y^2}$  symmetry whereas the tail sitting at the neighbouring oxygen sites are shaped according to  $\text{O-p}_x/\text{p}_y$  symmetries, indicating strong  $pd_\sigma$  bonding between Cu and O. Relatively large contribution of wannier orbitals on O- and V-sites indicate that the intrachain Cu1-Cu1 interaction is strongly mediated by the O atoms while interchain Cu1-Cu2 interaction is mediated by the V atoms respectively.  $pd\pi$  antibonds formed between  $\text{Cu-d}_{x^2-y^2}$  & O-p with somewhat stronger weights at O sites compared to that at V sites implies larger hybridization with O than V.

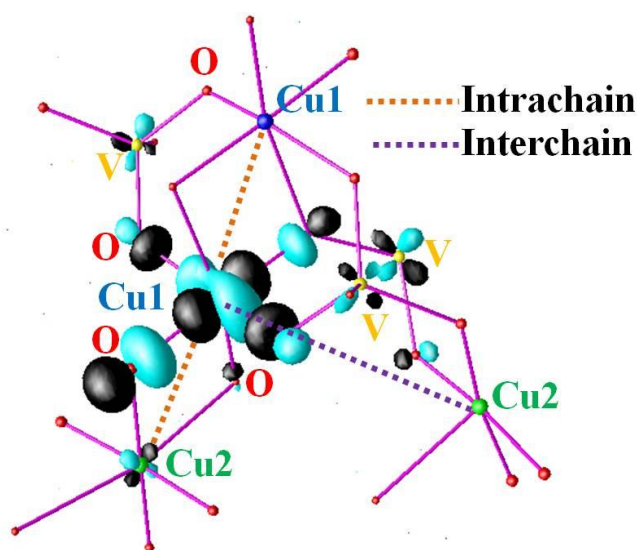
The paths for dominant magnetic interactions for  $\text{CuInO}(\text{VO}_4)$  is shown in Figure 5. We see that in  $\text{CuInO}(\text{VO}_4)$ , the intra-tetramer nearest neighbour interaction  $t_1$  is mediated by Cu1 d-O p-Cu1 d and Cu1 d-O p-V d-O p-Cu1 d super-exchange paths while  $t_2$  is mediated by Cu1 d-O p-Cu2 d super-exchange path. The inter-tetramer interactions  $t_3$  and  $t_4$  are mediated by Cu d-O p-V d-O p-Cu d and Cu d-O p-In s-O p-Cu d super super-exchange paths respectively.

Using superexchange formula  $4t^2/U$ , we can estimate the magnetic interactions between the magnetic ions from the values of hopping interaction with a choice of Hubbard parameter  $U$ . However, this provides

information of only the antiferromagnetic contributions. So, we used total energy calculation of different spin arrangements[14] and found out the magnetic exchange interactions by mapping the DFT energies to that of the Heisenberg model. Beyond GGA, the missing missing correlation energy at Cu sites is taken in terms of GGA+U calculations[15]. We have calculated the dominant magnetic exchanges of  $\text{CuInO}(\text{VO}_4)$  compound using  $U=6\text{eV}$ [16], which is shown in Table I. The positive sign of interactions indicates antiferromagnetic whereas negative sign implies ferromagnetic nature of exchange interactions.



**Figure 5:** The exchange paths for various magnetic interactions.



**Figure 6:** Wannier orbital for  $\text{CuInO}(\text{VO}_4)$  compound.

**TABLE 1:** The values of dominant hopping and magnetic exchange interactions in CuInO(VO<sub>4</sub>)

Interaction Type	Bond-distance (Å)	Path and Bond-angle	Hopping Intn. t's (meV)	$J=4t^2/U$ (meV)	Exchange Intn. J's (meV)
Intrachain	3.117	Cu1-O-Cu1; 89.76°	t1=150	J1=14.9	J1=22.8
	3.17	Cu1-O-Cu2; 88.22° & 107.55°	t2=136	J2=12.3	J2=7.7
Interchain	6.15	via In	t3=30	J3=0.6	J3= -0.06
	6.58	via V	t4=60	J4=2.4	J4=3.1

## 5. Conclusion

From microscopic investigations we conclude that the system appears weakly coupled Cu2-Cu1-Cu1-Cu2 AFM tetramer. The strongest intrachain interaction,  $J_1=22.8$  meV between Cu1-Cu1 is an order of magnitude stronger than that of Cu1-Cu2 intrachain interaction,  $J_2=7.7$  meV. Inter-tetramer interactions are  $J_3= -0.06$ meV and  $J_4=3.1$ meV which are weak compared to intra-tetramer interactions. We note that calculated small value of  $J_3$  is beyond the DFT accuracy limit, and the precise numerical value may not be trustworthy. Spin fluctuation prevent the long the long range order because the system is low dimensional. Therefore, Neel temperature ( $T_N$ ) should not be very high for this system.

## 6. Acknowledgements

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